

## **SSA and partonic intrinsic motion**

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## Outline

- Generalized pQCD approach with spin and  $\mathbf{k}_\perp$ -effects in distribution and fragmentation functions and elementary dynamics for  $pp \rightarrow hX$ ;
- SSA - theory: helicity formalism;
- SSA - phenomenology (1): Sivers and Collins effects vs. data;
- SSA - phenomenology (2):  $x_F < 0$ : the gluon Sivers function;
- SSA - phenomenology (3):  $D$  meson production and Drell-Yan processes;
- Conclusions and outlook

## Polarized cross sections: Helicity formalism

- ◊ **Helicity density matrices** within a  $\mathbf{k}_\perp$ -factorization scheme to describe parton spin states for a **polarized cross section**:

$$d\sigma^{A, \mathcal{S}_A + B, \mathcal{S}_B \rightarrow C + X} = \sum_{a,b,c,d,\{\lambda\}} \rho_{\lambda_a, \lambda'_a}^{a/A, \mathcal{S}_A} \hat{f}_{a/A, \mathcal{S}_A}(x_a, \mathbf{k}_{\perp a}) \otimes \rho_{\lambda_b, \lambda'_b}^{b/B, \mathcal{S}_B} \hat{f}_{b/B, \mathcal{S}_B}(x_b, \mathbf{k}_{\perp b}) \\ \otimes \hat{M}_{\lambda_c, \lambda_d; \lambda_a, \lambda_b} \hat{M}_{\lambda'_c, \lambda_d; \lambda'_a, \lambda'_b}^* \otimes \hat{D}_{\lambda_c, \lambda'_c}^{\lambda_C, \lambda'_C}(z, \mathbf{k}_{\perp C})$$

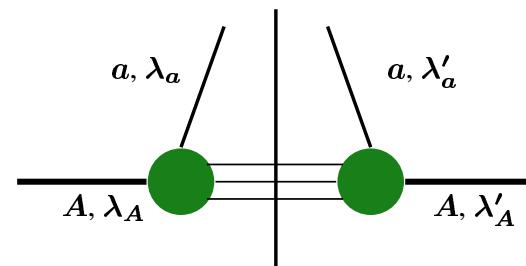
- $\rho_{\lambda_a, \lambda'_a}^{a/A, \mathcal{S}_A}$ : helicity density matrix of parton  $a$  inside hadron  $A$  with spin  $\mathcal{S}_A$
- $\frac{d\hat{\sigma}^{ab \rightarrow cd}}{dt} \simeq \sum_{\lambda_a, \lambda_b, \lambda_c, \lambda_d} |\hat{M}_{\lambda_c, \lambda_d; \lambda_a, \lambda_b}|^2$  (scattering amplitudes)
- $\hat{D}_{\lambda_c, \lambda'_c}^{\lambda_C, \lambda'_C} \equiv \sum_{X, \lambda_X} \hat{D}_{\lambda_C, \lambda_X; \lambda_c} \hat{D}_{\lambda'_C, \lambda_X; \lambda'_c}^*$   
i.e. the product of *fragmentation amplitudes* for  $c \rightarrow C + X$

## Helicity density matrices: phases from soft parts

- ◊ PDF at L.O. (and twist two) as the inclusive cross section for  $A \rightarrow a + X$ .  
Helicity density matrix of parton  $a \leftrightarrow$  to the hel. density matrix of hadron  $A$ .

$$\rho_{\lambda_a, \lambda'_a}^{a/A, S_A} \quad \hat{f}_{a/A, S_A}(x_a, \mathbf{k}_{\perp a}) = \sum_{\lambda_A, \lambda'_A} \rho_{\lambda_A, \lambda'_A}^{A, S_A} \hat{F}_{\lambda_A, \lambda'_A}^{\lambda_a, \lambda'_a}$$

$$\hat{F}_{\lambda_A, \lambda'_A}^{\lambda_a, \lambda'_a} \equiv \sum_{X_A, \lambda_X} \hat{\mathcal{F}}_{\lambda_a, \lambda_X; \lambda_A} \hat{\mathcal{F}}_{\lambda'_a, \lambda_X; \lambda'_A}^*$$



$$\begin{aligned} \hat{\mathcal{F}}_{\lambda_a, \lambda_X; \lambda_A}(x_a, \mathbf{k}_{\perp a}) &= \mathcal{F}_{\lambda_a, \lambda_X; \lambda_A}(x_a, k_{\perp a}) \exp[i\lambda_A \phi_a] \\ \Rightarrow \hat{F}_{\lambda_A, \lambda'_A}^{\lambda_a, \lambda'_a}(x_a, \mathbf{k}_{\perp a}) &= F_{\lambda_A, \lambda'_A}^{\lambda_a, \lambda'_a}(x_a, k_{\perp a}) \exp[i(\lambda_A - \lambda'_A) \phi_a] \end{aligned}$$

◊ same for FF (pion case):  $\hat{D}_{\lambda_c, \lambda'_c}^\pi = D_{\lambda_c, \lambda'_c}^\pi \exp[i(\lambda_c - \lambda'_c) \phi_\pi^H]$

## Partonic interpretation: notations

Spin and  $k_\perp$ -dependent quark distributions of leading twist (8):

$$\Delta \hat{f}_{s_x/\uparrow} \equiv \hat{f}_{s_x/\uparrow} - \hat{f}_{-s_x/\uparrow} = [F_{+-}^{+-} + F_{+-}^{-+}] \sin \phi_a \quad [h_{1T}, h_{1T}^\perp]$$

$$\Delta \hat{f}_{s_y/\uparrow} = -2 \operatorname{Im} F_{++}^{+-} + [F_{+-}^{+-} - F_{+-}^{-+}] \cos \phi_a \quad [h_1^\perp, h_{1T}, h_{1T}^\perp]$$

$$\Delta \hat{f}_{s_z/\uparrow} = 2 \operatorname{Re} F_{+-}^{++} \sin \phi_a \quad [g_{1T}^\perp]$$

$$\Delta \hat{f}_{s_x/+} = 2 \operatorname{Re} F_{++}^{+-} \quad [h_{1L}^\perp]$$

$$\Delta \hat{f}_{s_y/+} = \Delta \hat{f}_{s_y/A} = -2 \operatorname{Im} F_{++}^{+-} \quad [h_1^\perp]$$

$$\Delta \hat{f}_{s_z/+} = [F_{++}^{++} - F_{--}^{++}] \quad [g_1(\Delta f_a)]$$

$$\hat{f}_{a/\uparrow} = \hat{f}_{a/A} + \frac{1}{2} \Delta^N \hat{f}_{a/\uparrow} = [F_{++}^{++} + F_{--}^{++}] + 2 \operatorname{Im} F_{+-}^{++} \cos \phi_a .$$

$$[f_1, f_{1T}^\perp]$$

## Single Spin Asymmetries: $A^\dagger B \rightarrow \pi X$

- Configuration ( $AB \equiv pp$ ):

$AB$  center of mass frame -  $A$  along  $+Z$ -axis,  $\pi$  in the  $XZ$  plane with  $p_T$  parallel to  $+X$ -axis.  $S_A = \uparrow$  or  $\downarrow$  along  $Y$ -axis,  $S_B = 0$ .

- $k_\perp$ 's imply: non-planar  $ab \rightarrow cd$  scattering; different  $\perp$  directions.

Parton helicity amplitudes in  $pp$  c.o.m frame in terms of hel. amplitudes in partonic c.o.m. frame.

boost + rotations:  $\hat{M}_{\lambda_c, \lambda_d; \lambda_a, \lambda_b} = \hat{M}_{\lambda_c, \lambda_d; \lambda_a, \lambda_b}^0(\hat{s}, \hat{t}) e^{i\lambda_m \xi_m} e^{i(\lambda_a - \lambda_b)\phi_c''}$

Various combinations of “T-odd” effects in (un)polarized cross sections but

by including proper phases + numerical (8-dim., VEGAS) integration:

$$d\sigma^{\text{unp}} \simeq f_{a/p} \otimes f_{b/p} \otimes d\hat{\sigma} \otimes D_{\pi/c}$$

$$\begin{aligned} d\sigma^\uparrow - d\sigma^\downarrow &\simeq \Delta^N f_{a/p^\uparrow} \otimes f_{b/p} \otimes d\hat{\sigma} \otimes D_{\pi/c} && \text{“Sivers effect”} \\ &+ h_{1a} \otimes f_{b/p} \otimes d\Delta\hat{\sigma} \otimes \Delta^N D_{\pi/c^\uparrow} && \text{“Collins effect”} \end{aligned}$$

Notice: the explicit phases entering the convolutions are

$$\Delta^N f_{a/p^\uparrow} \otimes f_{b/p} \otimes d\hat{\sigma} \otimes D_{\pi/c} \rightarrow \sin(\phi_{S_A} - \phi_a) \equiv \cos \phi_a$$

$$h_{1a} \otimes f_{b/p} \otimes d\Delta\hat{\sigma} \otimes \Delta^N D_{\pi/c^\uparrow} \rightarrow \sin[\phi_{S_A} - (\phi_a + \phi_c'' - \xi_a - \tilde{\xi}_a + \xi_c + \tilde{\xi}_c + \phi_\pi^H)]$$

M. Anselmino, M. Boglione, U.D., E. Leader, F. Murgia PRD71 (05)

## Parameterization of the Sivers function

For each VALENCE quark  $q = u, d$  we adopt the factorized form

$$\Delta^N f_{q/p^\uparrow}(x, k_\perp) = 2 \mathcal{N}_q(x) h(k_\perp) f_{q/p}(x, k_\perp)$$

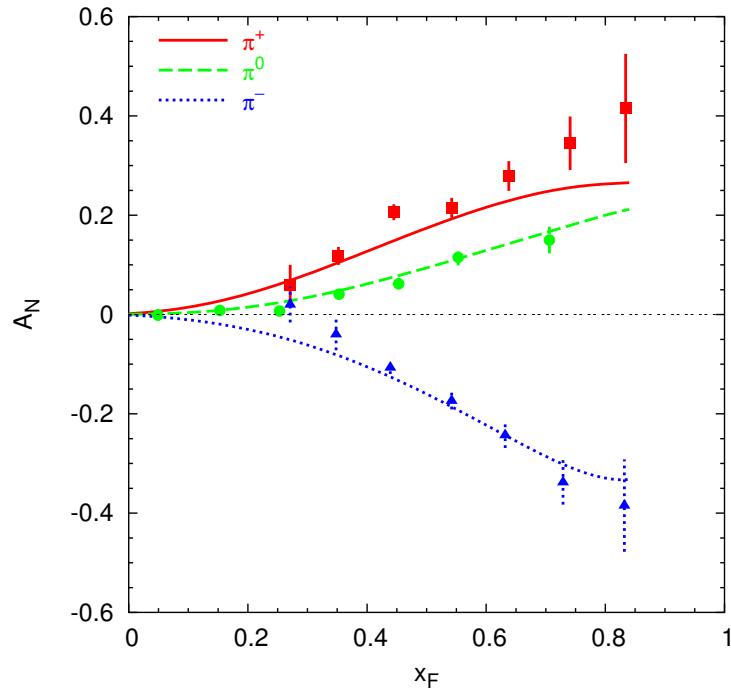
where

$$\mathcal{N}_q(x) = N_q x^{a_q} (1-x)^{b_q} \frac{(a_q + b_q)^{(a_q + b_q)}}{a_q^{a_q} b_q^{b_q}}$$

$$h(k_\perp) = \frac{2k_\perp M}{k_\perp^2 + M^2}$$

alternative  $k_\perp$  depend. :       $h'(k_\perp) = \sqrt{2e} \frac{k_\perp}{M'} e^{-k_\perp^2/M'^2}$

→ 7 parameters



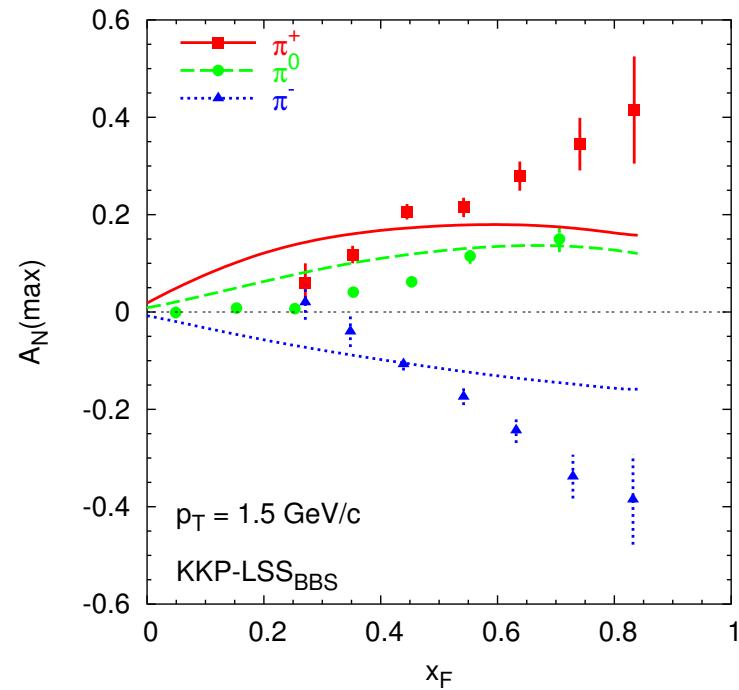
$A_N$  at  $E = 200$  GeV vs.  $x_F$  at  $p_T = 1.5$  GeV/c. Data are from [E704] PLB261-264 (1991).

Sivers effect [left] (valence-like).

$$N_u = +0.40 \quad a_u = 2.0 \quad b_u = 0.3$$

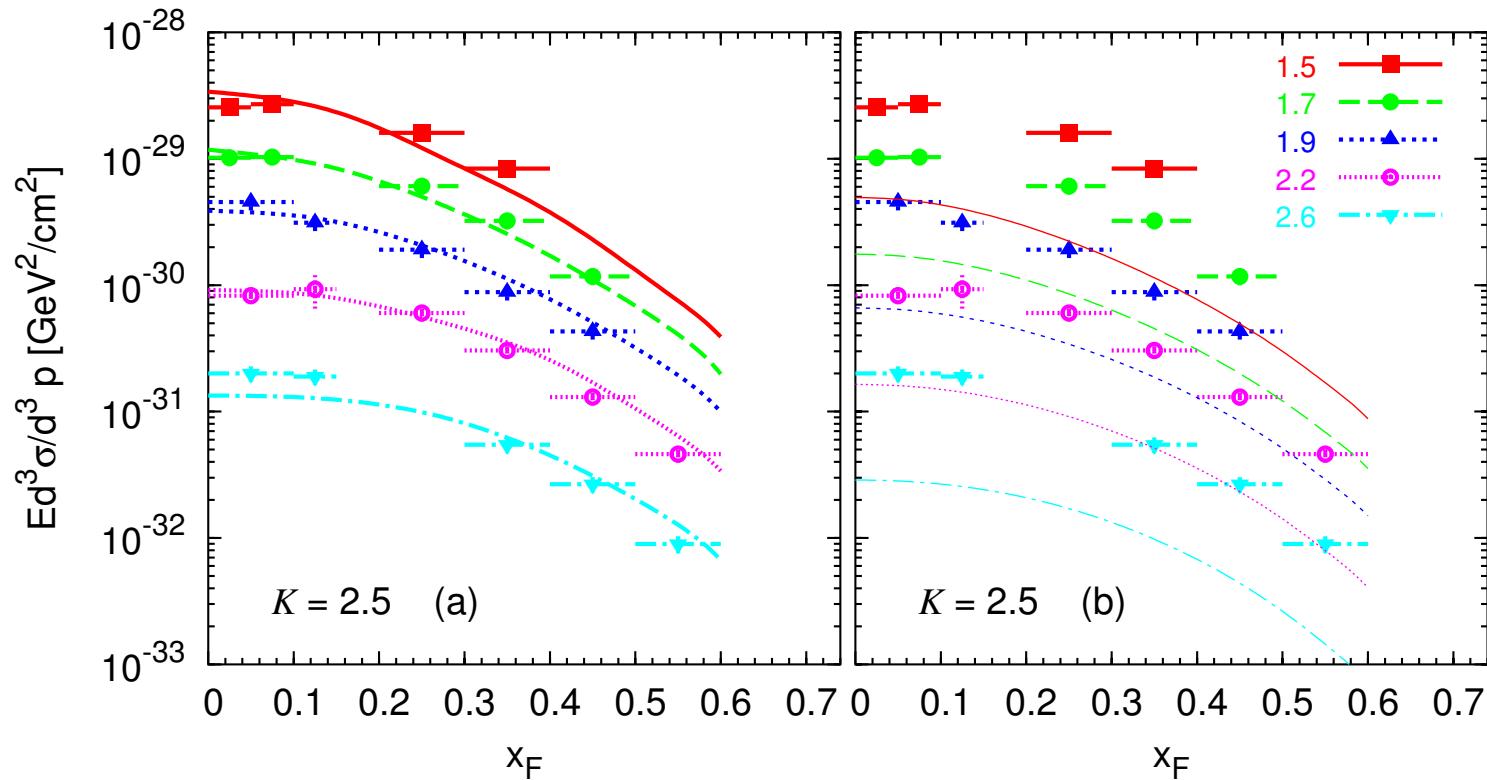
$$N_d = -0.90 \quad a_d = 2.0 \quad b_d = 0.2$$

$\simeq$  CONSISTENCY with SIDIS

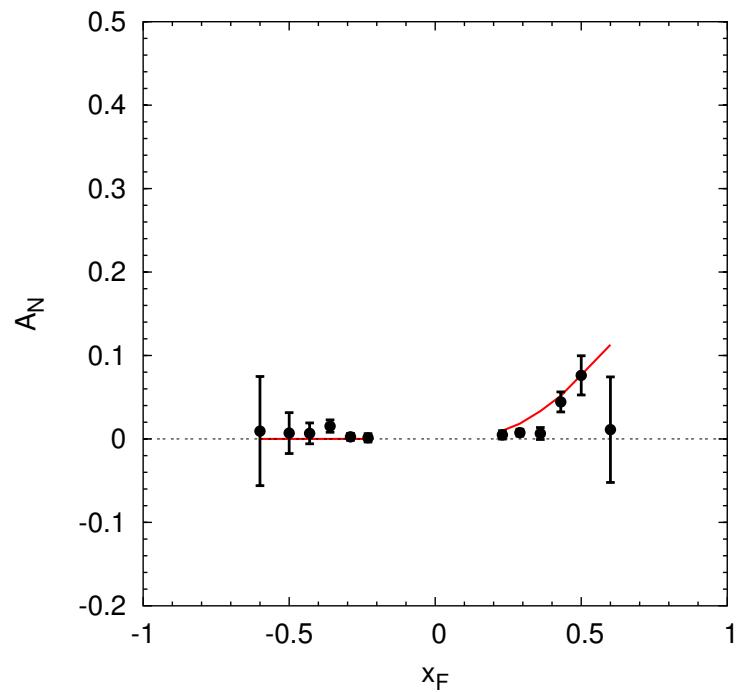
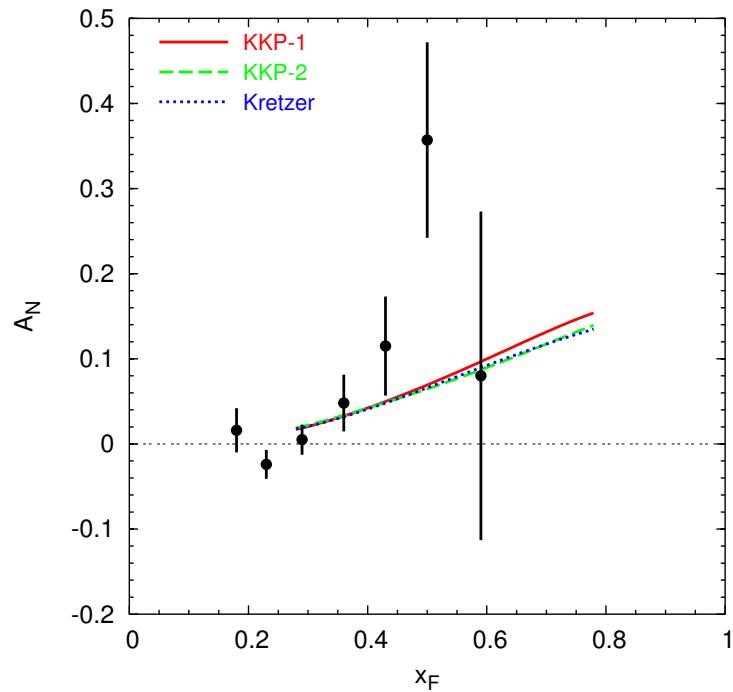


Collins effect [right](full saturated).

Transversity funct. and  
Collins funct. full saturated

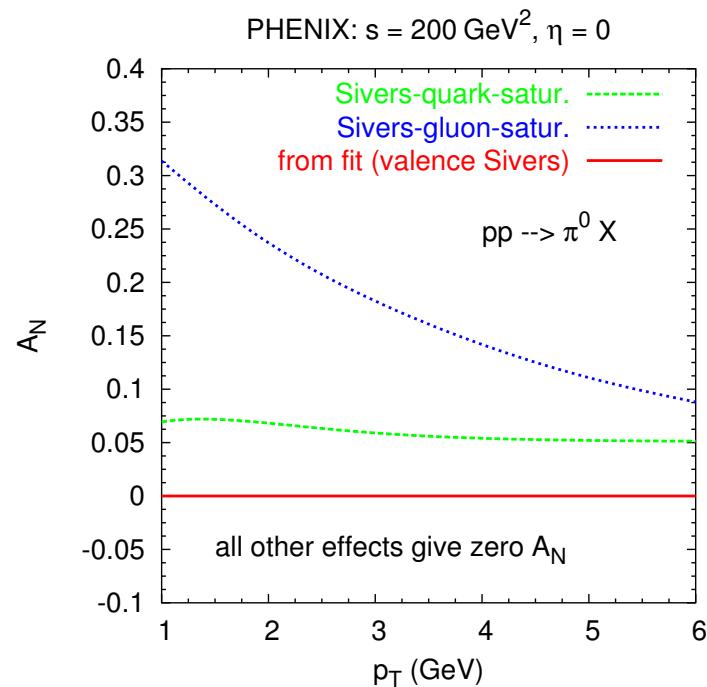
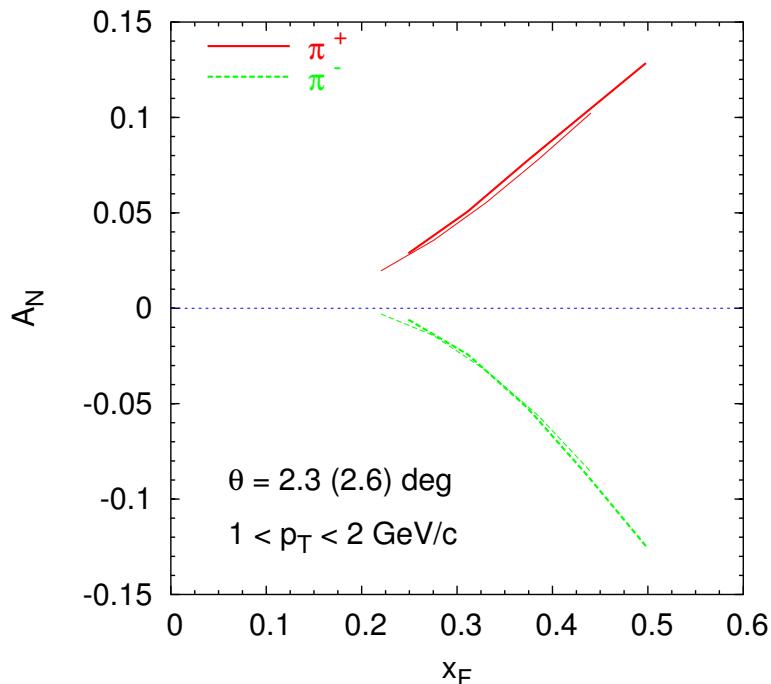


Estimates of  $\pi^0$  invariant cross sections **with (left)** and **without (right)**  $k_\perp$  effects at  $E = 200$  GeV vs.  $x_F$  for different  $p_T$  values. U.D. & F. Murgia PRD70 [05]. Data are from [BNL] Donaldson *et al.* PLB 73 (1978).



Predictions of  $A_N(pp \rightarrow \pi^0 X)$  in terms of Sivers effect alone [U.D. and F. Murgia PRD70 (05)] at  $\sqrt{s} = 200$  GeV and  $\eta = 3.8$  vs.  $x_F$ . Data are from [STAR] PRL92 (04).

Predictions of  $A_N(pp \rightarrow \pi^0 X)$  in terms of Sivers effect alone at  $\sqrt{s} = 200$  GeV and  $\eta = 4.1$  vs.  $x_F$ . STAR preliminary data.



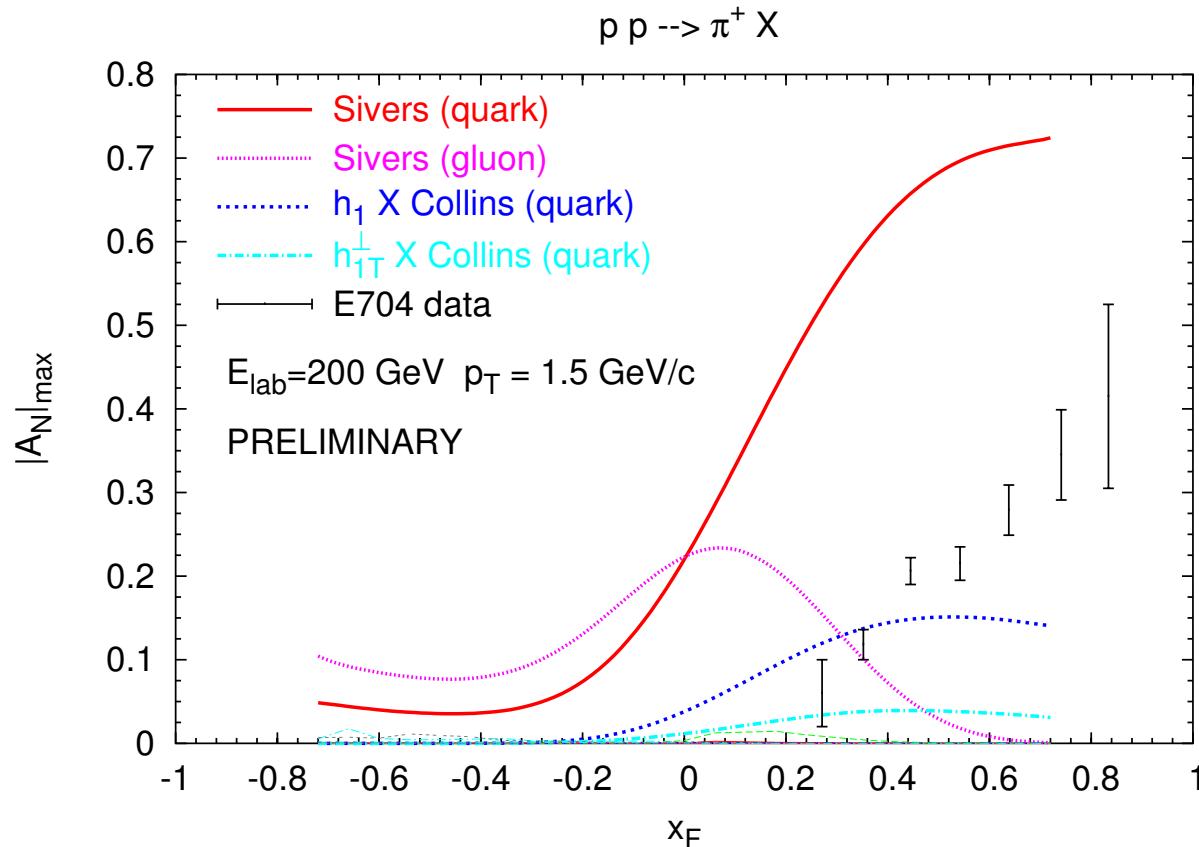
Predictions of  $A_N(pp \rightarrow \pi^\pm X)$  in terms of Sivers effect alone at  $\sqrt{s} = 200\ GeV$  and  $\theta = 2.3(2.6)\ deg.$  vs.  $x_F$  [BRAHMS kinematics].

Predictions of  $A_N(pp \rightarrow \pi^0 X)$  in terms of Sivers effect alone at  $\sqrt{s} = 200\ GeV$  and  $\eta = 0$  vs.  $p_T$  [PHENIX kinematics]. Consistent with preliminary data.  
Maximized effects are also shown.  
Collins effect suppressed!

## Consistency between Sivers functions from $p^\uparrow p \rightarrow \pi X$ and $\ell p^\uparrow \rightarrow \ell' \pi X$

- $\ell p \rightarrow \ell' \pi X$  (HERMES): moderately small  $x_B(x) \Rightarrow$  explore “sea region” and No gluons at LO;
- $pp \rightarrow \pi X$  (E704, STAR), large  $A_N$  at large  $x_F \Rightarrow$  constraint to the valence behaviour of TMD pdf’s. [U.D. and F. Murgia PRD70 (04)]  
gluon-sea cancelation at work at  $x_F \simeq 0$ ?
- Interplay of Collins effect ?  
New analysis [M.Anselmino et al. PRD71 (05)] with full microscopic dynamics  $\Rightarrow$  suppression of the Collins effect;

## All mechanisms at work in $p^\uparrow p \rightarrow \pi^+ X$



Contributions to  $A_N$  at  $\sqrt{s} = 20 \text{ GeV}$ , keeping only proper phases:  $\Delta f(x, k_\perp) = (2)f(x, k_\perp)$ .

## SSA at $x_F < 0$ and the gluon Sivers function

$p^\uparrow p \rightarrow \pi X$  in terms of  $\textcolor{red}{a}b \rightarrow cd$  [ $c$  fragments into  $\pi$ ]

$x_F > 0$  ( $x_1 > x_F$ )  $\rightarrow$  valence region of  $p^\uparrow$  and sea region for  $p$

$x_F < 0$  ( $x_1$  low)  $\rightarrow$  sea region of  $p^\uparrow$  and valence region of  $p$

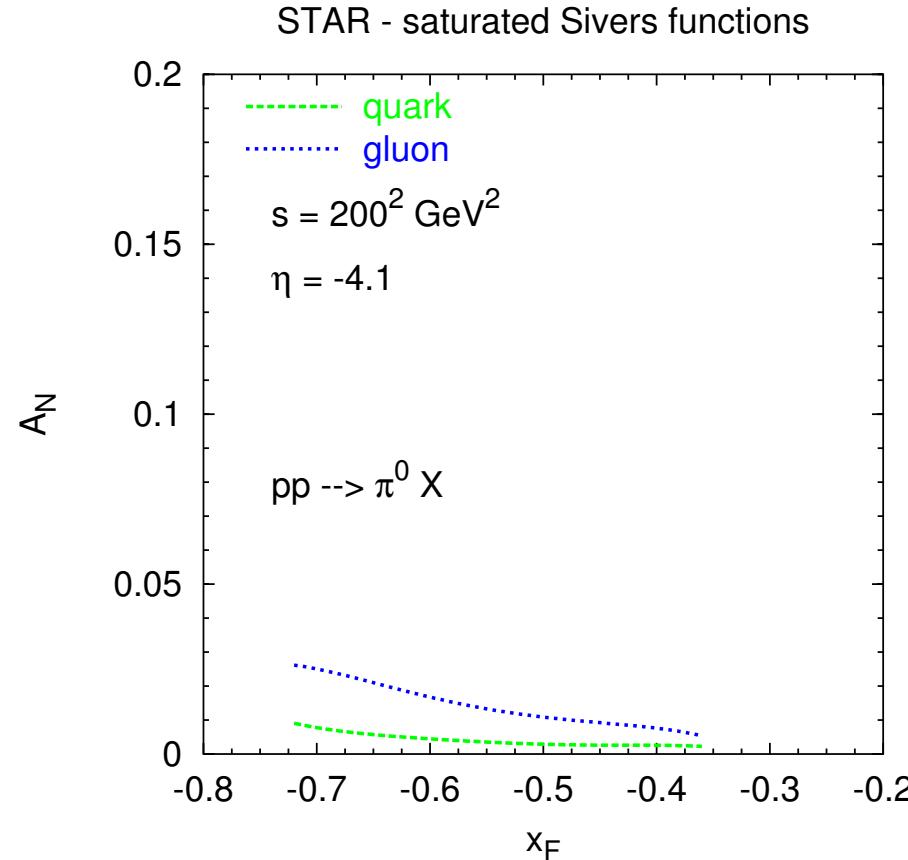
in particular

$x_F \rightarrow 1 \Rightarrow \hat{t} \rightarrow 0$  and  $\textcolor{red}{x}_1 \rightarrow 1$ :

- i)  $\hat{s}^2/\hat{t}^2$  terms dominate the partonic cross section
- ii)  $\textcolor{red}{q}b \rightarrow qb$  (larger PDF and FF) enhanced

$x_F \rightarrow -1 \Rightarrow \hat{u} \rightarrow 0$  and  $\textcolor{red}{x}_1 \rightarrow 0$ :

- i)  $\hat{s}^2/\hat{u}^2$  terms dominate the partonic cross section
- ii)  $\textcolor{red}{g}q \rightarrow qg$  (larger PDF and FF) enhanced



Estimates of  $A_N$  (Sivers effect), saturating the positivity bounds  $\left[ \Delta^N f_{a/p\uparrow}(x, k_\perp) = 2f_{a/p}(x, k_\perp) \right]$ , at  $\sqrt{s} = 200 \text{ GeV}$  and fixed  $\eta = 4.1$ . **5 times smaller than at  $\sqrt{s} = 20 \text{ GeV}!!!$**

## Phases and partonic dynamics [Sivers effect: $\cos \phi_a$ ]

partonic process:  $ab \rightarrow cd$

$$\hat{t} = (p_a - p_c)^2 = -2E_a E_c (1 - \cos \theta_{ac})$$

$$\hat{u} = (p_b - p_c)^2 = -2E_b E_c (1 - \cos \theta_{bc})$$

$$\mathbf{p}_c = E_c (\sin \theta_c, 0, \cos \theta_c) = E_c (\sin \theta_h, 0, \cos \theta_h) \quad [\mathbf{p}_c // \mathbf{p}_h]$$

$$\mathbf{p}_a = (k_{\perp a} \cos \phi_a, k_{\perp a} \sin \phi_a, E_a \cos \theta_a) \quad \mathbf{p}_b = (k_{\perp b} \cos \phi_b, k_{\perp b} \sin \phi_b, E_b \cos \theta_b)$$

with

$$\cos \theta_a \simeq 1 + \mathcal{O}(k_{\perp}^2/s) \quad \cos \theta_b \simeq -1 + \mathcal{O}(k_{\perp}^2/s)$$

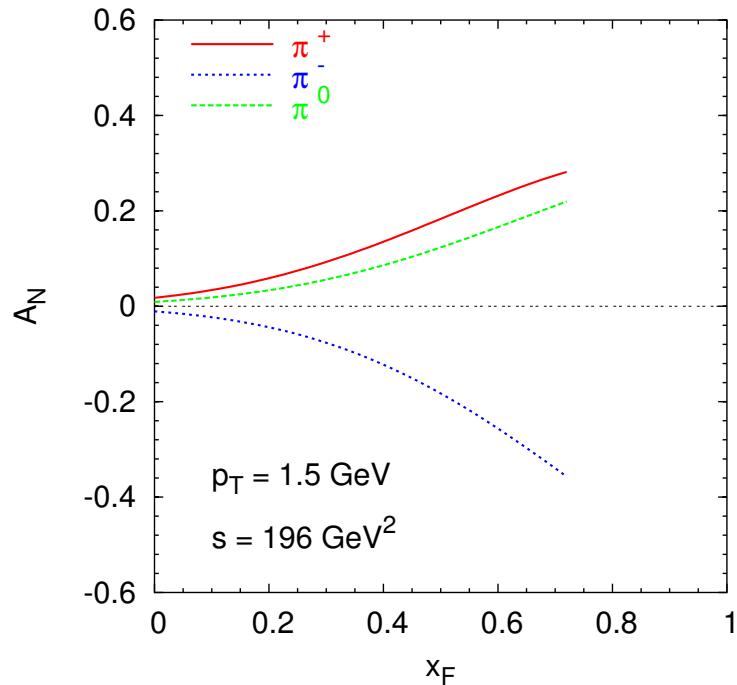
$$\cos \theta_{ac} = k_{\perp a}/E_a \cos \phi_a \sin \theta_h + \cos \theta_h$$

$$\cos \theta_{bc} = k_{\perp b}/E_b \cos \phi_b \sin \theta_h - \cos \theta_h$$

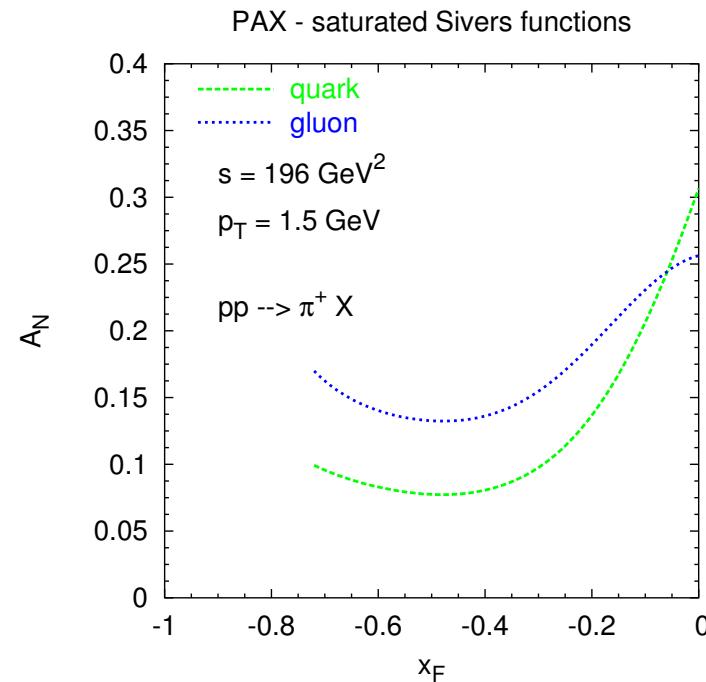
$\theta_h \rightarrow 0 (x_F > 0)$ :  $\hat{t} \rightarrow 0$  enforced for  $\phi_a \rightarrow 0$  almost independent of  $\phi_b$

$\theta_h \rightarrow \pi (x_F < 0)$ :  $\hat{u} \rightarrow 0$  enforced for  $\phi_b \rightarrow 0$  almost independent of  $\phi_a$

- at STAR  $|\eta| \approx 4$  i.e.  $\theta_h \approx 2(178)$  degrees  
⇒ at  $x_F < 0$  (all range)  $\hat{u}$  dominance and almost no  $\hat{t}$  contamination  
⇒  $d\hat{\sigma}$  almost independent of  $\phi_a$ : strong suppression from the integration of the  $\cos \phi_a$  Sivers factor;  
 $x_F < 0$  data seem not able to constrain the gluon Sivers pdf
- at E704,  $\sqrt{s} = 20$  GeV with  $p_T > 1 - 1.5$  GeV  
⇒  $10(20) < \theta_h < 170(160)$  degrees (depending on  $x_F$ )  
⇒ also at  $x_F \rightarrow -1$  still a contamination from  $\hat{t}$ , i.e. dependence on  $\phi_a$ , in  $d\hat{\sigma}$  and therefore a less severe suppression;  
but No data
- at PAX  $A_N$  in the negative  $x_F$  region: a potential tool.



Predictions of  $A_N(pp \rightarrow \pi X)$  in terms of Sivers effect alone at  $\sqrt{s} = 14$  GeV and  $p_T = 1.5$  GeV vs.  $x_F$  [PAX kinematics].



Estimates of  $A_N(pp \rightarrow \pi^+ X)$  in terms of Sivers effect alone, saturating the positivity bounds, at  $\sqrt{s} = 14$  GeV and  $p_T = 1.5$  GeV vs.  $x_F$  [PAX kinematics].

## Gluon Sivers function from $p^\uparrow p \rightarrow DX$ at RHIC

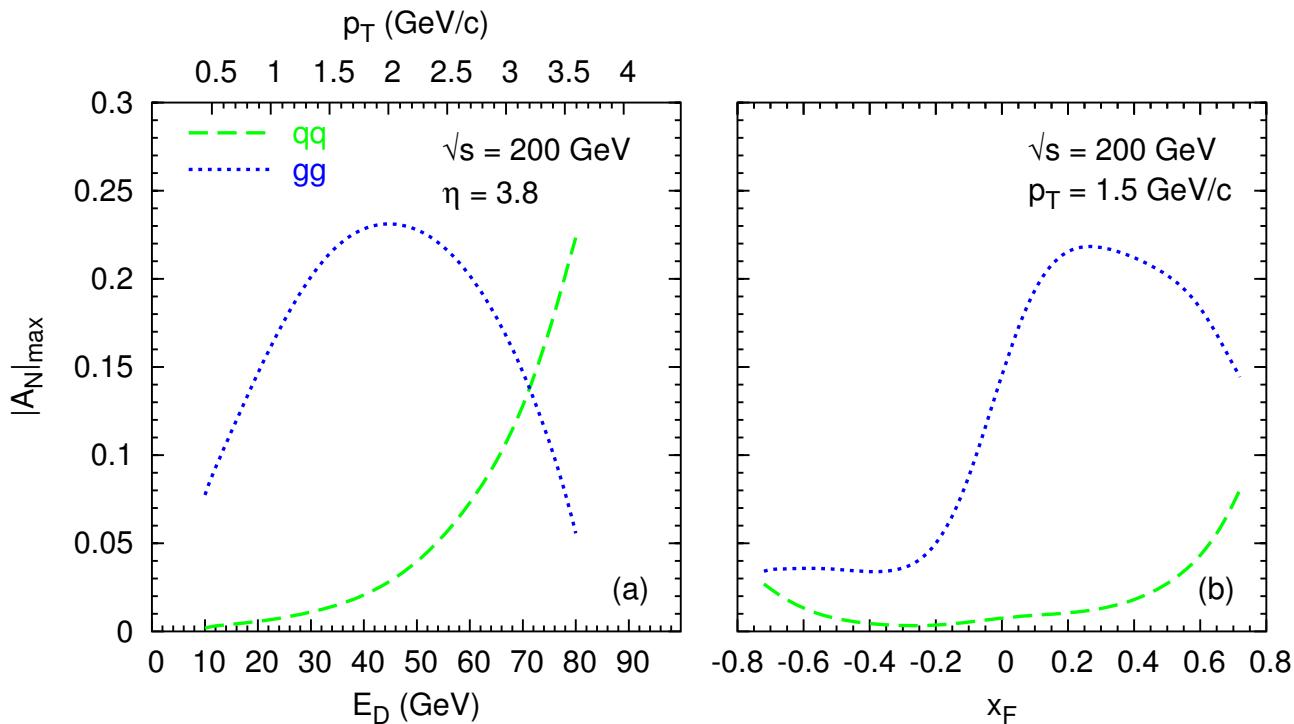
M. Anselmino, M. Boglione, U.D., E. Leader, F. Murgia PRD70 (04)

- unpol. cross-sections:  $q\bar{q} \rightarrow c\bar{c}$  +  $gg \rightarrow c\bar{c}$  (dominant: up to 10 times)
- helicity formalism with full  $\mathbf{k}_\perp$ :  $d\sigma^\uparrow - d\sigma^\downarrow$  develops various “T-odd” combinations

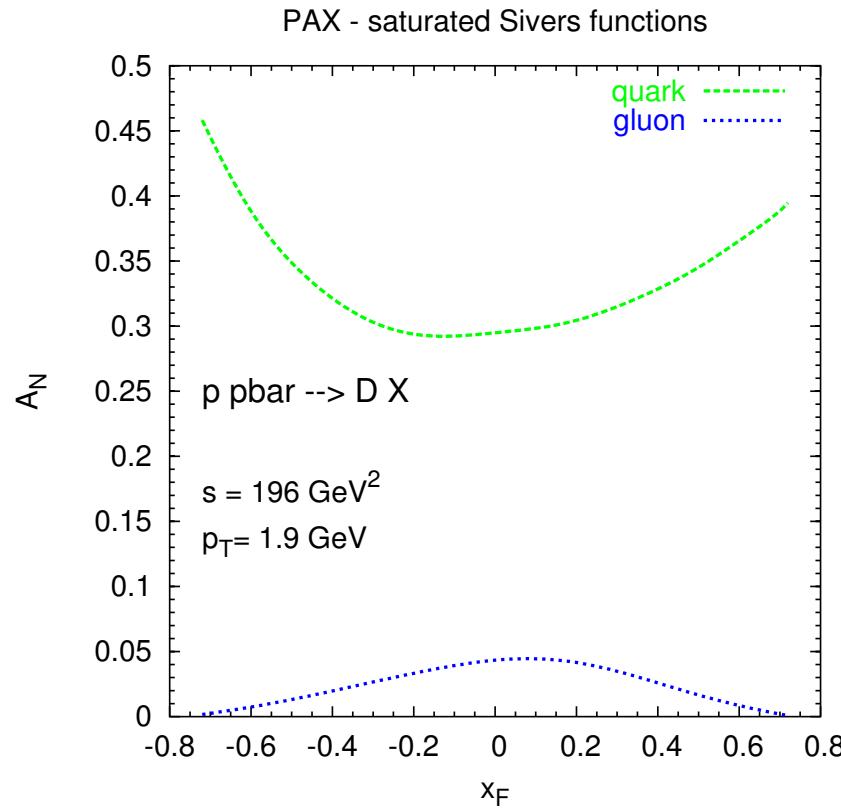
NO  $h_1 \Delta^N D_{D/c^\uparrow}$  (zero transverse single spin transfer)

- proper phases: integration washes out all terms BUT Sivers effect
- gluon dominance up to  $x_F = 0.5$  ( $x$ -saturated):  
 $A_N \neq 0 \Rightarrow$  gluon Sivers effect

Also from jet correlations in  $p^\uparrow p$  reactions (D. Boer, W. Vogelsang PRD69 [04]).



Estimates for  $A_N$  in terms of Sivers effect,  $x_F$  – saturated from quarks or gluons at  $\sqrt{s} = 200$  GeV vs.  $x_F$  at fixed  $\eta = 3.8$  (left) and at fixed  $p_T = 1.5$  GeV/c (right).  
FF set: Cacciari *et al.* PRD55 (1997).



Estimates of  $A_N$  for  $p^\uparrow p \rightarrow DX$  at PAX, collider mode,  $\sqrt{s} = 14 \text{ GeV}$ , at fixed  $p_T = 1.9 \text{ GeV}$ . **Access to the quark Sivers function. NO Collins effect!!**

## Access to the Sivers effect and universality: SSA in $p^\uparrow p \rightarrow \ell^+ \ell^- X$ processes

[M. Anselmino, U.D., F. Murgia, PRD67 (03)]

Differential cross sections in the variables:  $M^2, y, \mathbf{q}_T$     [ $q_T^2 \ll M^2$ ]

Angular distribution of the lepton pair production plane: integrated over.

$$d\sigma^\uparrow - d\sigma^\downarrow \simeq \Delta^N f_{q/A^\uparrow}(x_a, \mathbf{k}_{\perp a}) \otimes \hat{f}_{\bar{q}/B}(x_b, \mathbf{k}_{\perp b}) d\hat{\sigma}^{q\bar{q} \rightarrow \ell^+ \ell^-}$$

Other mechanisms for SSA in Drell-Yan processes:

$$\sum_q h_1(x_a, \mathbf{k}_{\perp a}) \otimes \Delta^N f_{\bar{q}^\uparrow/B}(x_b, \mathbf{k}_{\perp b}) \otimes d\Delta\hat{\sigma}^{q\bar{q} \rightarrow \ell^+ \ell^-} \quad [\text{D. Boer PRD60 (99)}]$$

$$d\Delta\hat{\sigma} = d\hat{\sigma}^{\uparrow\uparrow} - d\hat{\sigma}^{\uparrow\downarrow} \simeq \cos 2\phi \quad [\phi = \widehat{P_A N}_{\ell^+ \ell^-}] \quad \int d\phi \Rightarrow 0$$

Universality [J. Collins PLB536 (02)]:

$$\Delta^N f_{q/A^\uparrow}(x, k_\perp)|_{DY} = -\Delta^N f_{q/A^\uparrow}(x, k_\perp)|_{DIS} ??? \dots$$

## Conclusions and outlook

- A generalized pQCD approach to SSA and unpol. cross sections:  
LO-pQCD + spin and new  $k_\perp$ -dependent PDF and FF;
- $A_N(p^\uparrow p \rightarrow \pi X)$ : Complete  $k_\perp$ -helicity formalism + phenomenology.  
*True detailed microscopic dynamics (proper phases)  $\Rightarrow$*   
Sivers effect results **confirmed**;  
Collins effect **suppressed** (against former claims!);
- good agreement with recent STAR data on  $A_N$  both at  $x_F > 0$  and  $x_F < 0$  in terms of Sivers effect alone and valence like behaviour.
- $A_N(p^\uparrow p \rightarrow \pi X)$  at  $x_F < 0$ : strong suppression (of all mechanisms!!)  
by phase conspiracy;
  - at large energies (RHIC) and  $\theta \rightarrow \pi$  all contributions, even saturated, give almost zero SSA.
  - at intermediate energies (PAX) the  $x_F < 0$  region may constrain the gluon Sivers function;

- $A_N(p^\dagger p \rightarrow \pi X)$  at  $x_F = 0$  vs.  $p_T$  [PHENIX]: strong suppression; ONLY SIVERS effect (gluon and/or sea quark) at work!
- SSA with  **$D$  meson production** as a tool to extract the **gluon Sivers function (RHIC)** and the **quark Sivers function (PAX)**.
- SSA in **Drell-Yan**: a tool to extract the **quark Sivers function** and to test the universality condition on TMD pdf.
- Combined analysis of different processes and in different kinematical configurations to map the TMD structure of hadrons.